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Numerical Study of Hydrogen Crossover Evolution Inside the Proton Exchange Membrane Fuel Cell under Dynamic Load

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ABSTRACT: Hydrogen (H₂) crossover in proton exchange membrane fuel cells (PEMFCs) reduces performance and poses safety risks, but its behavior under rapidly changing loads, which are common in vehicles, is not well understood. To address this, we developed a three-dimensional, two-phase, non-isothermal model that tracks H₂ from dissolution in the anode, through transport across the membrane, to reaction at the cathode. The analysis shows that diffusion dominates whereas convection contributes little. Key findings are as follows: H₂ crossover reduces the open-circuit voltage by 210 mV and raises cathode temperature by approximately 0.2°C; reducing the membrane thickness from 20 to 5 μm increases the crossover current density fourfold (from 2.8–3.6 to 11.4–13.2 mA cm⁻²); under rapid load changes, transient undershoots of 0.8–1.72 mA cm⁻² occur because the H₂ concentration drops quickly whereas water and thermal conditions adjust slowly; and a variation of approximately 1 mA cm⁻² along the flow channel indicates that local H₂ distribution and membrane hydration strongly affect transport. Overall, H₂ crossover under dynamic loads is governed by diffusion as modified by local water and heat distribution, with significant differences between channel and rib regions. These results help predict and mitigate fuel cell degradation in practical applications.

KEYWORDS: Proton exchange membrane fuel cells; hydrogen crossover; numerical simulation; three-dimensional model; dynamic load

1 Introduction

Proton exchange membrane fuel cells (PEMFCs) are considered to be promising power sources for vehicles due to their high efficiency, zero-emission, and quiet operation [1,2]. However, unlike steady-state laboratory testing, automotive applications subject PEMFCs to severe dynamic operating conditions—including rapid load transients during acceleration/deceleration and frequent start-stop cycles—that result in dramatic changes in local water, thermal and chemical conditions [3], fundamentally altering H₂ transport phenomena [4]. Under these dynamic conditions, the undesired hydrogen (H₂) crossover exhibits transient behaviors significantly distinct from steady-state diffusion, yet its spatiotemporal evolution remains poorly understood despite causing membrane thinning and degradation [5–7]. When H₂ crossover occurs, hydrogen is transported to the cathode electrode, where it undergoes reactions generating heat and water, but without contributing to electrical output. Furthermore, hydrogen peroxide radicals can form at potentials below 0.628 V [8]. The elevated local temperature and the presence of hydrogen peroxide radicals can initiate thermal and chemical degradation of the perfluorosulfonic acid (PFSA) ionomer, potentially leading to membrane deterioration, pinhole development [9] and associated safety risks such as hydrogen

explosions [5,10,11] over long-term operation. While these long-term degradation mechanisms and safety hazards are not explicitly modeled in the present work, understanding the transient H₂ crossover dynamics and local thermal responses addressed herein provides an essential basis for predicting these phenomena. Therefore, it is crucial to evaluate local H₂ crossover inside PEMFCs, especially in dynamic operating scenarios, considering both efficiency and safety.

Despite extensive research on H₂ crossover, significant knowledge gaps persist regarding its behavior under realistic dynamic operating conditions. Current understanding relies heavily on steady-state investigations, creating a critical disconnect with automotive applications characterized by rapid load transients. Experimental studies have employed both direct and indirect methodologies to quantify H₂ crossover [12]. Direct techniques such as gas chromatography [13], mass spectrometry [10], and volumetric methods [14] provide accurate permeation data, while indirect electrochemical methods including linear sweep voltammetry (LSV) [15], cyclic voltammetry (CV) [16], potential step method (PSM) [17], and galvanostatic charging method (GCM) [18,19] offer operational convenience. However, most electrochemical techniques integrate crossover over the entire membrane area, obscuring local spatial distributions. Although specialized segmented plates [20] and multi-port bipolar plates [21] have enabled localized measurements, such studies remain limited to steady-state conditions with nitrogen-purged cathodes to eliminate reaction interference. Consequently, the online monitoring of local H₂ crossover during actual fuel cell discharge under dynamic loads remains technically unachievable due to intricate mass transfer and reaction complexities.

Theoretical modeling offers an alternative to resolve spatial and temporal distributions, providing equipment-independent, cost-effective, and accessible benefits [4,22–25], yet current approaches exhibit systematic simplifications. Most existing models treat H₂ crossover as steady-state Fickian diffusion through the membrane [26–29], neglecting the complete physicochemical process encompassing gas dissolution, pressure-driven convection, interfacial release, and electrochemical oxidation at the cathode [14,30,31]. While some studies acknowledge pressure-driven convection [32,33] and cathode reaction kinetics [23], few models integrate these mechanisms within a comprehensive multi-physics framework. More critically, the transition from steady-state to dynamic operation has received limited attention [8,34], despite H₂ permeation fluctuations under transient conditions directly governing localized degradation rates. These collective limitations reveal a fundamental research gap: the absence of comprehensive three-dimensional transient models capable of resolving local H₂ crossover dynamics and their coupling with thermal and electrochemical responses under realistic automotive operating scenarios.

To contextualize these limitations and clarify the distinct contributions of the present work, [Table 1](#) systematically compares the key modeling features (dimensionality, thermal coupling, transient capability, and H₂ crossover physics) of the most relevant prior studies [23,27,28] against the proposed framework.

Table 1: Model parameters.

Feature	Vilekar and Datta [23]	Nguyen and Fushinobu [28]	Ji et al. [27]	Present Work
Dimensionality	1D/Simplified	2D	2D/Lumped-parameter	3D full-field
Two-phase	No	Yes	Yes	Yes
Thermal coupling	Isothermal (implicit)	Not considered	Simplified/Isothermal	Non-isothermal (full energy equation)
H ₂ crossover physics	Steady-state diffusion only	Steady gas permeation	Fick's law coupling	Dissolution-diffusion-convection-reaction

(Continued)

Table 1 (continued)

Feature	Vilekar and Datta [23]	Nguyen and Fushinobu [28]	Ji et al. [27]	Present Work
Transient capability	Steady-state	Steady-state	Quasi-steady	Fully transient
Dynamic load support	Not supported	Not supported	Not supported	Step/ramp load transients

Note: Vilekar and Datta [23] is a baseline steady-state dissolution-diffusion model; Nguyen and Fushinobu [28] represents 2D two-phase steady-state gas crossover; Ji et al. [27] employs Fick's law coupling but lacks full transient thermal-mass coupling.

The main contributions of this study are summarized as follows:

(1) Development of a coupled multi-physics transient model.

Unlike previous studies simplifying H_2 crossover as steady-state diffusion, we establish a three-dimensional, two-phase, non-isothermal transient model integrating H_2 dissolution, diffusion-convection transport, and electrochemical oxidation, coupled with local mass, charge, and heat transfer.

(2) Experimental validation under transient conditions.

The model is validated against experimental data for both steady-state polarization and dynamic voltage responses (step load changes), addressing the critical gap of transient H_2 crossover measurement that direct experimental techniques cannot easily achieve online.

(3) Discovery of transient undershoot mechanisms and parametric effects.

We reveal the undershoot phenomenon ($0.8\text{--}1.72\text{ mA cm}^{-2}$) driven by the competition between concentration drop and membrane hydration lag, and quantify effects of membrane thickness ($5\text{--}20\text{ }\mu\text{m}$), stoichiometry ($0.9\text{--}2.0$), and humidity ($20\%\text{--}80\%\text{ RH}$).

(4) Revealing spatial-temporal variations and establishing foundation for degradation prediction.

We identify contrasting channel-rib evolution patterns during load transients (channel-rib differences of $0.1\text{--}0.7\text{ mA cm}^{-2}$) and quantify humidity-induced spatial variations ($\sim 0.8\text{ mA cm}^{-2}$ increase), establishing a mechanistic foundation for subsequent coupling with chemical degradation models.

The remainder of this paper is organized as follows. [Section 2](#) presents the model assumptions. [Section 3](#) describes the development of the PEMFC model coupled with H_2 crossover, including process features and geometric domain, governing equations, reaction kinetics, and boundary conditions. [Section 4](#) presents model implementation and verification, followed by results and discussion of parametric studies and design implications. Finally, [Section 5](#) summarizes the principal findings, discusses model limitations, and outlines future work.

2 Model Assumption

The model is based on the following reasonable assumptions [16]:

- (1) All gases behave as ideal gases.
- (2) Product water is initially produced in the gas phase and may undergo phase change during transport through the porous layers, forming liquid water.
- (3) The membrane is defect-free, and H_2 crossover flux varies spatially with temperature, membrane water content, and current density.
- (4) Permeated H_2 reacts electrochemically with O_2 .

3 PEMFC Model Coupled with H₂ Crossover

3.1 Process Features and Geometric Domain

As illustrated in Fig. 1a, the operation of hydrogen-powered fuel cells is characterized by frequent fluctuations in current loading over time. This results in real-time variations in internal heat and mass transfer, as well as fluctuations in electrochemical reaction rates within the cell. At the anode catalyst layer (aCL), hydrogen (H₂) undergoes the hydrogen oxidation reaction (HOR), generating protons and electrons at platinum (Pt) catalyst sites. Protons are transported to the cathode catalyst layer (cCL) through ionomer and proton exchange membrane (PEM) phases, while electrons travel through porous carbon materials and external circuitry to the cathode side, performing electrical work during this process. At the cathode, oxygen (O₂) reacts with protons and electrons to form water, releasing significant heat. The transport of mass, heat, and charge within the fuel cell is depicted in Fig. 1b.

The PEM serves the crucial functions of conducting protons and separating anode fuel from cathode oxidant. However, H₂ crossover is unavoidable due to the porous nature of membranes and tends to increase with PEM aging [16,35]. As shown in Fig. 1c, a fraction of H₂ permeates through the membrane via diffusion and convective transport to reach the cCL [33]. Upon reaching the cathode, permeated H₂ undergoes chemical or electrochemical reactions. Most studies indicate that H₂ is completely consumed by O₂ through electrochemical reactions, primarily due to high cathode electrode potentials [15,23,36,37]. Therefore, as shown in Fig. 1c, H₂ reaching the cathode is dissociated into electrons and protons through HOR, which then combine with O₂ within the cCL to undergo oxygen reduction reaction (ORR), thereby forming a local fuel cell. Although HOR and ORR in this local fuel cell consume hydrogen and oxygen to produce water, they do not generate external electrical work. As a result, enthalpy changes from both half-reactions are released as heat, potentially leading to local hot spot formation [38].

H₂ crossover processes are significantly affected by fuel cell operating conditions; therefore, a three-dimensional (3D) PEMFC model was utilized in this study to investigate local H₂ crossover behavior. The model consists of a single channel with a length of 10 mm, as illustrated in Fig. 1b. The geometric domain encompasses the bipolar plate (BP), gas channel (GC), gas diffusion backing layer (GDB), microporous layer (MPL), catalyst layer (CL), and PEM, where lowercase letters “a” and “c” represent the anode and cathode sides, respectively. The geometric configuration and cross-sectional parameters are consistent with specifications detailed in our previous publications on two-dimensional modeling [39,40]. Additional geometric parameters for this model are listed in Table 2.

3.2 Model Equations

This study investigates the evolution of H₂ crossover in PEMFCs under dynamic loading conditions using a three-dimensional, two-phase, non-isothermal, transient model. The model is based on multi-physics conservation equations governing mixture gas flow, gas species transport, dissolved H₂ transport, liquid water transport, membrane water content, energy, and charge. The two-phase flow, heat, and charge transport components are primarily adapted from the one-dimensional and two-dimensional transient models of multi-physics coupled phenomena in PEMFCs developed by Wang et al. [39,41]. Additionally, dissolved H₂ transport in the electrolyte is modeled based on the steady-state model developed by Vilekar and Datta [23].

The symbols and nomenclature used in Eqs. (1)–(37) are defined in Table 3.

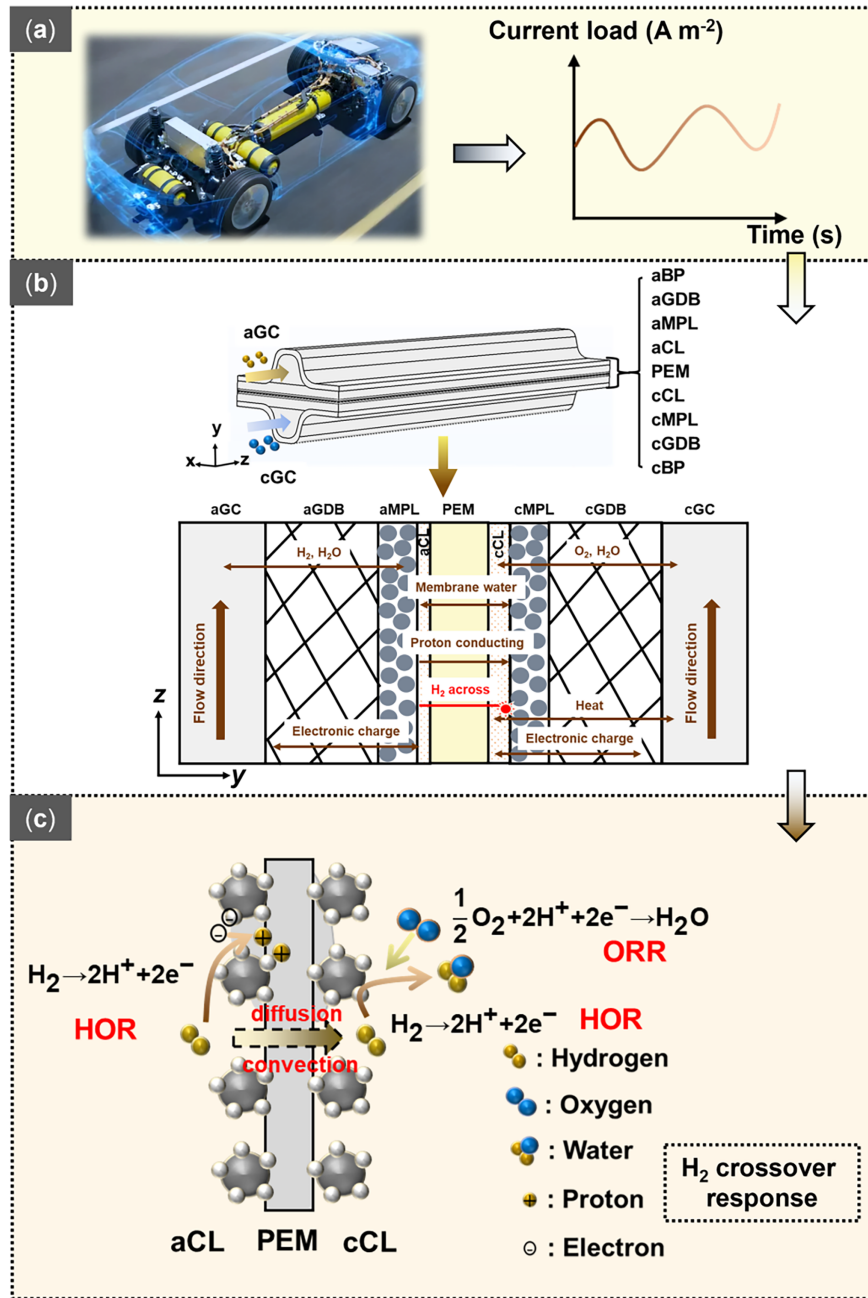


Figure 1: Schematics of (a) dynamic load of fuel cell power system, (b) single-channel fuel cell 3D geometric model and internal transport process, and (c) detailed mechanism of H₂ crossover (dissolution at anode, transport across membrane, and electrochemical reaction at cathode).

Table 2: Model parameters.

Gas Channel Length, Height, Bottom Width, Top Width (mm)	10, 0.7, 1, 0.5
Land length, width, thickness (mm)	10, 0.75, 0.1
Thickness of aCLs and cCL (μm)	4, 8
Thickness of MPLs (μm)	30
Thickness of GDBs (μm)	190
Thickness of PEM (μm)	15
Porosity of CLs, MPLs, and GDBs	0.3, 0.6, 0.75
Volume fraction of the Ionomer in CLs	0.3
Hydraulic permeability of liquid water through the membrane (m^2)	1.6×10^{-18}
Mass transport coefficient for gas H_2 to dissolved H_2 (s^{-1})	100,000
Tortuosity of the ionomer	1.5
Membrane water dynamic viscosity (Pa s)	4.7×10^{-4}
Dry membrane density (kg m^{-3})	2000
Enthalpy change for ORR (J mol^{-1})	-241,826
Standard entropy change for ORR ($\text{J mol}^{-1} \text{K}^{-1}$)	-44.43
Exchange current density for ORR (mA cm^{-2})	1.5×10^{-6}
Exchange current density for HOR (mA cm^{-2})	5
Contact angle of CL, MPL, GDB ($^\circ$)	100, 130, 140
Surface tension coefficient (N m^{-1})	0.0625
Water evaporation rate (s^{-1})	100
Water condensation rate (s^{-1})	100

Table 3: Nomenclature.

Symbol	Description	Unit
a_{Pt}	Electrocatalytic surface area per unit volume	m^2/m^3
A_{m}	Area of membrane-electrode assembly parallel to flow channel	m^2
A_{ch}	Cross-sectional area of flow channel	m^2
C_i	Concentration of gas species i	mol/m^3
C_{H_2}	Hydrogen concentration	mol/m^3
$C_{\text{H}_2,\text{diss}}$	Concentration of dissolved H_2	mol/m^3
$C_{\text{H}_2,\text{eq}}$	Equilibrium concentration of saturated dissolved H_2	mol/m^3
$C_{\text{H}_2}^{\text{ref}}$	Reference hydrogen concentration	mol/m^3
C_{m}	Concentration of membrane water	mol/m^3
C_{O_2}	Oxygen concentration	mol/m^3
$C_{\text{O}_2}^{\text{ref}}$	Reference oxygen concentration	mol/m^3
$C_{\text{p,g}}$	Specific heat capacity of gas phase	$\text{J}/(\text{kg K})$
$C_{\text{p,i}}$	Specific heat capacity of species i	$\text{J}/(\text{kg K})$
$C_{\text{p,l}}$	Specific heat capacity of liquid water	$\text{J}/(\text{kg K})$
$C_{\text{p,s}}$	Specific heat capacity of solid phase	$\text{J}/(\text{kg K})$
D_i^{eff}	Effective diffusion coefficient of gas species i	m^2/s
$D_{\text{H}_2,\text{diss}}$	Effective diffusion coefficient of dissolved H_2 in electrolyte	m^2/s

(Continued)

Table 3 (continued)

Symbol	Description	Unit
$D_{H_2,w}$	Diffusion coefficient of dissolved H_2 in pure water	m^2/s
D_m^{eff}	Effective diffusion coefficient of membrane water	m^2/s
F	Faraday constant	C/mol
f_v	Water volume fraction within ionomer	–
$I(t)$	Time-varying current load	A/m^2
$i_{0,a}^{ref}$	Reference exchange current density for HOR	A/m^2
$i_{0,c}^{ref}$	Reference exchange current density for ORR	A/m^2
J_a	Volumetric current density of HOR at anode	A/m^3
J_c	Volumetric current density of ORR at cathode	A/m^3
J_{cross}	Volumetric current density from HOR of permeated H_2 at cathode	A/m^3
k^{eff}	Effective thermal conductivity	$W/(m K)$
k_{diss}	Mass transport coefficient for H_2 dissolution	s^{-1}
K_{ml}	Hydraulic permeability of liquid water through membrane	m^2
M_{H_2}	Molar mass of hydrogen	kg/mol
M_{H_2O}	Molar mass of water	kg/mol
M_{O_2}	Molar mass of oxygen	kg/mol
n_d	Electro-osmotic drag coefficient	–
p_{ml}	Membrane water pressure	Pa
$p_{g,a}$	Gas pressure at anode inlet	Pa
$p_{g,c}$	Gas pressure at cathode inlet	Pa
p_{H_2}	Partial pressure of H_2 in pore space	Pa
p_{sat}	Saturation vapor pressure	Pa
R	Universal gas constant	$J/(mol K)$
RH_a, RH_c	Relative humidity at anode/cathode inlets	–
s	Liquid water saturation	–
S_e	Source term for electrons	A/m^3
S_g	Mass source term for gas mixture	$kg/(m^3 s)$
S_H	Source term for protons	A/m^3
S_{H_2}	Source term for hydrogen gas	$mol/(m^3 s)$
$S_{H_2,diss}$	Source term for H_2 dissolution from pores into ionomer	$mol/(m^3 s)$
S_{H_2O}	Source term for water vapor	$mol/(m^3 s)$
$S_{H_2,w}$	Solubility of hydrogen in pure water	$mol/(m^3 Pa)$
S_i	Source term for gas species i	$mol/(m^3 s)$
S_l	Mass source term for liquid water	$kg/(m^3 s)$
S_{ml}	Mass source for liquid mixture in membrane	$kg/(m^3 s)$
S_{O_2}	Source term for oxygen gas	$mol/(m^3 s)$
$S_{r,diss}$	Mass source for dissolved H_2 consumption by HOR in CCL	$mol/(m^3 s)$
S_T	Heat source term	W/m^3
S_T^{v-d}	Heat associated with water adsorption	W/m^3
S_T^{v-l}	Heat associated with water condensation	W/m^3
S_{v-d}	Quantity of gaseous water adsorbed by ionomers	$mol/(m^3 s)$

(Continued)

Table 3 (continued)

Symbol	Description	Unit
S_{v-1}	Amount of gaseous water condensing into liquid water	mol/(m ³ s)
T	Temperature	K
u_a	Inlet flow velocity at anode	m/s
u_c	Inlet flow velocity at cathode	m/s
u_g	Velocity of gas mixture	m/s
$u_{H_2,diss}$	Convection velocity of liquid mixture (membrane water + dissolved H ₂)	m/s
u_l	Velocity of liquid water	m/s
ε	Volumetric porosity of porous materials	–
ε_m	Ionomer volume fraction	–
η_a	Overpotential for HOR at anode	V
η_c	Overpotential for ORR at cathode	V
η_{cross}	Overpotential for HOR at cathode	V
ξ_a	Stoichiometric ratio of H ₂ at anode	–
ξ_c	Stoichiometric ratio of O ₂ at cathode	–
ρ_g	Density of gas mixture	kg/m ³
ρ_l	Density of liquid water	kg/m ³
ρ_s	Density of solid phase	kg/m ³
σ_e^{eff}	Effective electron conductivity	S/m
σ_m^{eff}	Effective proton conductivity	S/m
τ	Tortuosity of ionomer	–
ϕ_e	Electronic potential	V
$\phi_{eq,a}$	Equilibrium potential for HOR	V
$\phi_{eq,c}$	Equilibrium potential for ORR	V
ϕ_m	Proton potential	V
ΔH_{hor}	Enthalpy change of HOR	J/mol
ΔH_{orr}	Enthalpy change of ORR	J/mol
μ_l	Dynamic viscosity of water	Pa s

3.2.1 H₂ Crossover Transport

The conservation of H₂ in the electrolyte involves three key processes: (1) dissolution of H₂ at the anode from pores into water-saturated ionomer [37], (2) diffusion and convection of molecular H₂ through the membrane from ACL to CCL, driven by concentration and pressure gradients [42], and (3) hydrogen oxidation reaction of dissolved H₂ at the cathode [23], as well as release of dissolved H₂ from the electrolyte into pores within the CCL, as illustrated in Fig. 2.

These processes are mathematically described by the dissolved hydrogen transport equation group (Eqs. (1)–(10)).

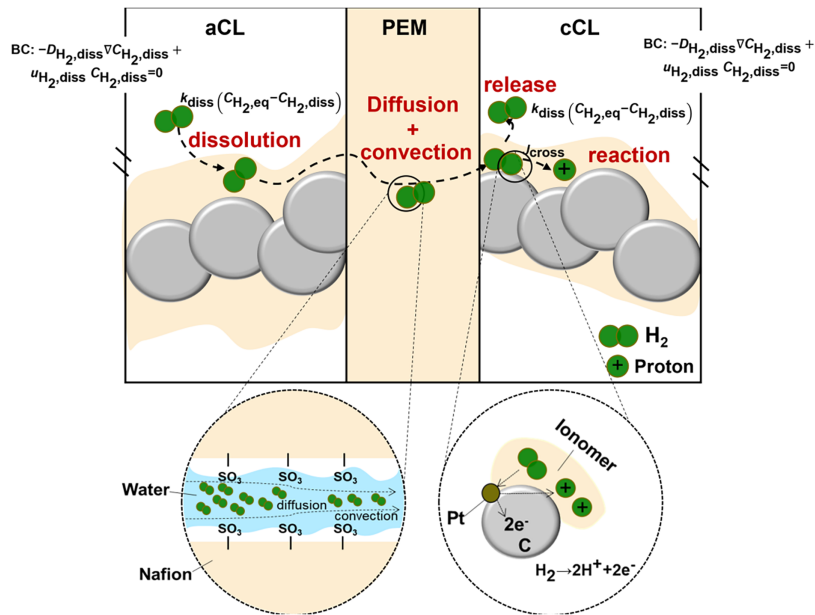


Figure 2: Schematic illustration of H₂ crossover mechanism showing the three key processes: (1) dissolution of H₂ into ionomer at aCL, (2) diffusion and convection through PEM, and (3) electrochemical reaction at cCL, with detailed insets of Nafion microstructure and Pt catalyst reaction.

Dissolved H₂ transport in ionomer phases of the CL and PEM is described by the following mass conservation equation:

$$\frac{\partial C_{H_2,diss}}{\partial t} + \nabla \cdot (C_{H_2,diss} \mathbf{u}_{H_2,diss}) + \nabla \cdot (-D_{H_2,diss} \nabla C_{H_2,diss}) = S_{H_2,diss} - S_{r,diss} \quad (1)$$

where $C_{H_2,diss}$ is the concentration of dissolved H₂; $\mathbf{u}_{H_2,diss}$ denotes the convection velocity of the liquid mixture (membrane water and dissolved H₂) in the membrane; $D_{H_2,diss}$ represents the effective diffusion coefficient of dissolved H₂ in the electrolyte. The effect of membrane water content on H₂ crossover is reflected in the water-dependent diffusion coefficient, which is calculated using the following equation [42–44]:

$$D_{H_2,diss} = \frac{\varepsilon_m f_v}{\tau} D_{H_2,w} \quad (2)$$

where ε_m , f_v and τ represent the ionomer volume fraction, water volume fraction within the ionomer, and tortuosity of the ionomer, respectively. $D_{H_2,w}$ is the diffusion coefficient of dissolved H₂ in pure water, which can be obtained from the following equation [43]:

$$D_{H_2,w} = 4.9e^{-6} \exp\left(\frac{-16,560.885}{RT}\right) \quad (3)$$

Convection of the liquid mixture follows Darcy’s law and is analyzed using the mass conservation equation for the liquid mixture (Eq. (4)) and Darcy’s equation (Eq. (5)), applicable to the CLs and PEM [44,45]. Given the negligible concentration of dissolved H₂, physical properties of the mixed solution, including density and viscosity, are approximated by those of liquid water.

$$\frac{\partial}{\partial t}(\varepsilon_m f_v \rho_l) + \nabla \cdot \rho_l \mathbf{u}_{\text{H}_2, \text{diss}} = S_{\text{ml}} \quad (4)$$

$$\mathbf{u}_{\text{H}_2, \text{diss}} = -f_v \frac{K_{\text{ml}}}{\mu_l} \nabla p_{\text{ml}} \quad (5)$$

where ε_m is the volume fraction of ionomer; ρ_l is the density of liquid water; K_{ml} is the hydraulic permeability of liquid water through the membrane; μ_l is the dynamic viscosity of water; and p_{ml} denotes membrane water pressure.

To quantitatively evaluate the relative contributions of diffusion and convection to hydrogen transport across the membrane, an order-of-magnitude analysis is conducted. The diffusion flux is expressed as $J_{\text{diff}} = -D_{\text{H}_2, \text{diss}} \cdot \frac{\Delta C_{\text{H}_2, \text{diss}}}{\delta_{\text{mem}}}$, where $\Delta C_{\text{H}_2, \text{diss}}$ is the concentration difference across the membrane and δ_{mem} denotes the membrane thickness. The convective flux is estimated by $J_{\text{conv}} = C_{\text{H}_2, \text{diss}} \cdot f_v \cdot \frac{K_{\text{ml}}}{\mu_l} \cdot \frac{\Delta p_{\text{ml}}}{\delta_{\text{mem}}}$, with Δp_{ml} representing the pressure difference. Using typical parameter values for a PEMFC— $D_{\text{H}_2, \text{diss}} \sim 10^{-10} \text{ m}^2 \text{ s}^{-1}$, $\Delta C_{\text{H}_2, \text{diss}} \sim 100 \text{ mol m}^{-3}$, $\delta_{\text{mem}} \sim 10^{-5} \text{ m}$, $C_{\text{H}_2, \text{diss}} \sim 10^{-1} \text{ mol m}^{-3}$, $f_v \sim 10^{-1}$, $K_{\text{ml}} \sim 10^{-19} \text{ m}^2$, $\mu_l \sim 10^{-4} \text{ kg m}^{-1} \text{ s}^{-1}$, and $\Delta p_{\text{ml}} \sim 10^4 \text{ Pa}$ —the fluxes are estimated to be $J_{\text{diff}} \sim 10^{-5} \text{ mol m}^{-2} \text{ s}^{-1}$ and $J_{\text{conv}} \sim 10^{-8} \text{ mol m}^{-2} \text{ s}^{-1}$. Under the simulated operating conditions, diffusion dominates over convection by approximately three orders of magnitude. Although convection is negligible under the conditions considered, the full convection–diffusion formulation is retained for two main reasons: (i) it enhances the model's generality, enabling its applicability to scenarios with more significant pressure gradients, such as during transient startup/shutdown, variations in membrane compression, or high differential pressure operation and (ii) it facilitates a more comprehensive mechanistic analysis of crossover dynamics without imposing a priori assumptions.

The term $S_{\text{H}_2, \text{diss}}$ in Eq. (1) represents the source term for H_2 dissolution from pores into water-saturated ionomer in the CLs, expressed as follows [43]:

$$S_{\text{H}_2, \text{diss}} = k_{\text{diss}} (C_{\text{H}_2, \text{eq}} - C_{\text{H}_2, \text{diss}}) \quad (6)$$

where k_{diss} represents the mass transport coefficient, and $C_{\text{H}_2, \text{eq}}$ denotes the equilibrium concentration of saturated dissolved H_2 , as described by Henry's law [43,46].

$$C_{\text{H}_2, \text{eq}} = S_{\text{H}_2, \text{w}} p_{\text{H}_2} \quad (7)$$

where p_{H_2} represents the partial pressure of H_2 in the pore space, and $S_{\text{H}_2, \text{w}}$ denotes the solubility of hydrogen in pure water, calculated using the following equation:

$$S_{\text{H}_2, \text{w}} = 0.54 \exp \left(-48.1611 + \frac{5528.45}{T} + 16.8893 \ln \left(\frac{T}{100} \right) \right) \quad (8)$$

The term $S_{\text{r}, \text{diss}}$ in Eq. (1) represents the mass source for dissolved H_2 consumption due to HOR in the CCL, calculated using the following equation [43]:

$$S_{\text{r}, \text{diss}} = \frac{J_{\text{cross}}}{2F} \quad (9)$$

where J_{cross} denotes the volumetric current density generated by HOR in the CCL, and F is the Faraday constant.

The term S_{ml} in Eq. (4) represents the mass source for the liquid mixture in the membrane, accounting for water adsorption from the gas phase to the membrane phase, H_2 dissolution, and dissolved H_2 consumption by HOR processes, calculated using the following equations:

$$S_{ml} = \begin{cases} S_{v-d}M_{H_2O} + (S_{H_2,diss} - S_{r,diss})M_{H_2}(cCL) \\ S_{v-d}M_{H_2O} + S_{H_2,diss}M_{H_2}(aCL) \end{cases} \quad (10)$$

where S_{v-d} denotes the quantity of gaseous water adsorbed by ionomers; calculation of this term follows equations presented in our previous work [41].

This concludes the dissolved hydrogen equation group. The resulting crossover current density J_{cross} (Eq. (9)) feeds into the charge transport (Section 3.2.5) and energy transport (Section 3.2.6) equation groups as a source term.

3.2.2 Mixture Gas Transport

This subsection presents the gas-phase mixture equation group (Eqs. (11)–(13)), governing the overall mass conservation of the gas mixture in gas channels and porous media.

The mass conservation equation of gas mixture in GCs, GDLs, MPLs, and CLs is solved as follows [39]:

$$\frac{\partial}{\partial t}(\varepsilon(1-s)\rho_g) + \nabla \cdot \rho_g \mathbf{u}_g = S_g \quad (11)$$

where ρ_g denotes the density of the gas mixture; \mathbf{u}_g is the velocity of the gas mixture, obtained using Darcy's law [39,47]; S_g represents the mass source; ε is the volumetric porosity of porous materials; and s is the liquid water saturation.

In Eq. (11), the source terms S_g denoted as represent mass consumption and generation during phase transitions and electrochemical reactions. Within the ACL, these terms account for H_2 consumption through HOR, H_2 dissolution from the gas phase into the ionomer, as well as water condensation and adsorption. The calculations are detailed as follows:

$$S_g = -\frac{J_a}{2F}M_{H_2} - M_{H_2O}S_{v-1} - M_{H_2O}S_{v-d} - M_{H_2}S_{H_2,diss} \quad (12)$$

where J_a represents the volumetric current density of HOR at the anode; M is the molar mass; and S_{v-1} indicates the amount of gaseous water condensing into liquid water in pores, as detailed in our previous work [41].

Within the CCL, these terms account for water generation and O_2 consumption by ORR, water phase transitions between gas-liquid-membrane phases, and H_2 dissolution from pores into ionomer. The calculations are shown as follows:

$$S_g = \frac{J_c}{2F}M_{H_2O} - \frac{J_c}{4F}M_{O_2} - M_{H_2O}S_{v-1} - M_{H_2O}S_{v-d} - M_{H_2}S_{H_2,diss} \quad (13)$$

where J_c is the volumetric current density at the cathode.

This concludes the mixture gas equation group. The mass source terms S_g couple this group to the electrochemical reaction kinetics (Section 3.3) and the dissolved hydrogen transport (Section 3.2.1).

3.2.3 Gas Species Transport

This section presents the gas species transport equation group (Eqs. (14)–(19)), tracking individual components (H_2 , O_2 , N_2 , and H_2O) via convection-diffusion equations.

The mass conservation of gas species in GCs, GDLs, MPLs, and CLs is solved as follows [41]:

$$\frac{\partial}{\partial t}(\varepsilon(1-s)C_i) + \nabla \cdot (C_i \mathbf{u}_g) + \nabla \cdot (-D_i^{\text{eff}} \nabla C_i) = S_i \quad (14)$$

where C_i is the gas species concentration; D_i^{eff} denotes the effective diffusion coefficient, calculated as detailed in our previous work [41]; and S_i represents the source terms due to electrochemical reactions and phase transitions, with values varying by species. At the anode, species i can be H_2 or H_2O , while at the cathode, species may include O_2 , H_2 , N_2 , and H_2O .

The source terms for H_2 and H_2O species in the ACL account for H_2 consumption by HOR, H_2 dissolution into ionomer, and water phase transitions, as outlined below.

$$S_{H_2} = -\frac{J_a}{2F} - S_{H_2,\text{diss}} \quad (15)$$

$$S_{H_2O} = -S_{v-1} - S_{v-d} \quad (16)$$

The source terms for O_2 , H_2 , and H_2O in the CCL account for O_2 consumption by ORR, H_2 dissolution into ionomer, H_2O generation by ORR, and water phase transitions between gas-liquid-membrane phases, as given below.

$$S_{O_2} = -\frac{J_c}{4F} \quad (17)$$

$$S_{H_2} = -S_{H_2,\text{diss}} \quad (18)$$

$$S_{H_2O} = \frac{J_c}{2F} - S_{v-1} - S_{v-d} \quad (19)$$

This concludes the species transport equation group. The species fluxes provide boundary conditions for the dissolved hydrogen equation group (Section 3.2.1) via Henry's law.

3.2.4 Liquid and Membrane Water Transport

This section comprises the water transport equation group (Eqs. (20) and (21)), addressing liquid water in porous media and membrane water governed by diffusion and electro-osmotic drag.

The mass conservation of liquid water in GDLs, MPLs, and CLs is governed by the following equation [48]:

$$\frac{\partial}{\partial t}(\varepsilon s \rho_l) + \nabla \cdot \rho_l \mathbf{u}_l = S_l \quad (20)$$

where ρ_l represents the density of liquid water; \mathbf{u}_l denotes the velocity of liquid water, determined by Darcy's law [39,47]; and S_l represents the mass source term resulting from water condensation, calculated as $M_{H_2O} S_{v-1}$.

The transport of membrane water in CLs and PEM involves two key mechanisms: membrane water diffusion driven by concentration gradients and electro-osmotic drag driven by potential gradients. The mass conservation equation is given as follows [41]:

$$\frac{\partial}{\partial t}(\epsilon_m C_m) + \nabla \cdot (-D_m^{\text{eff}} \nabla C_m - \frac{n_d}{F} \sigma_m^{\text{eff}} \nabla \phi_m) = S_{v-d} \quad (21)$$

where C_m is the concentration of membrane water in the membrane; D_m^{eff} is the effective diffusion coefficient; n_d is the electro-osmotic drag coefficient; σ_m^{eff} is the effective proton conductivity; ϕ_m is the proton potential; and S_{v-d} represents the mass source term for water adsorption by the ionomer. The methods for calculating these physical and transport parameters have been detailed in our previous work [41].

This concludes the water transport equation group. The membrane water content (characterized by the water volume fraction f_v in Eq. (2)) directly modifies the diffusion coefficient in the dissolved hydrogen equation group.

3.2.5 Charge Transport

This section presents the charge transport equation group (Eqs. (22)–(25)), describing proton and electron conduction via Ohm's law.

The charge balance of electrons in BPs, GDLs, MPLs, and CLs, and protons in CLs and PEM, follows Ohm's law and is determined using the following equation [41].

$$\nabla \cdot (-\sigma_m^{\text{eff}} \nabla \phi_m) = S_H \quad (22)$$

$$\nabla \cdot (-\sigma_e^{\text{eff}} \nabla \phi_e) = S_e \quad (23)$$

where σ_e^{eff} denotes the effective electron conductivity; ϕ_e is the electronic potential; and S_H and S_e are the source terms for protons and electrons, respectively.

The proton source in the aCL is derived from HOR, which generates protons. In contrast, the proton source in cCL includes proton consumption during ORR and proton production from permeated H_2 via HOR. The calculation equation is given as follows:

$$S_H = \begin{cases} J_a & \text{(aCL)} \\ -J_c + J_{\text{cross}} & \text{(cCL)} \end{cases} \quad (24)$$

The electron source in the aCL is attributed to electron generation from HOR. In contrast, the electron source in the cCL encompasses both electron consumption during ORR and electron generation from HOR due to H_2 crossover.

$$S_e = \begin{cases} -J_a & \text{(aCL)} \\ J_c - J_{\text{cross}} & \text{(cCL)} \end{cases} \quad (25)$$

This concludes the charge transport equation group. The current densities J_a and J_c serve as source terms for the gas-phase and energy equation groups.

3.2.6 Energy Transfer

This section presents the energy transport equation group (Eqs. (26) and (27)), governing heat conduction, convection, and generation within the cell, including the reversible and irreversible heat released by hydrogen crossover at the cathode.

The energy conservation in GDLs, MPLs, and CLs is addressed using the following equation, which incorporates energy transport mechanisms such as heat conduction, diffusion, and convection [41].

$$\frac{\partial}{\partial t} \left((1-\varepsilon) \rho_s C_{p,s} T + \varepsilon s \rho_l C_{p,l} T + \varepsilon (1-s) \rho_g C_{p,g} T \right) + \nabla \cdot (-k^{\text{eff}} \nabla T) + \nabla \cdot \left(\sum_i (-D_i^{\text{eff}} \nabla C_i) M_i C_{p,i} T \right) + \nabla \cdot (\rho_l \mathbf{u}_l C_{p,l} T + \rho_g \mathbf{u}_g C_{p,g} T) = S_T \quad (26)$$

where ρ denotes density, C_p represents specific heat capacity, and T is temperature. The term k^{eff} refers to effective thermal conductivity, while subscripts s, g, and l denote solid, gaseous, and liquid phases, respectively. Additionally, S_T represents the heat source, accounting for heat release associated with phase transitions, proton and electron conduction, and electrochemical reactions in CLs. Notably, the newly introduced reversible and irreversible heat source terms caused by HOR due to H_2 crossover in the CCL are included, as shown in the following equation:

$$S_T = \begin{cases} \sigma_e^{\text{eff}} + \sigma_m^{\text{eff}} \nabla \phi_m^2 + J_a \left(\phi_{\text{eq},a} - \frac{-\Delta H_{\text{hor}}}{2F} \right) + J_a \eta_a + S_T^{v-1} + S_T^{v-d} & \text{(aCL)} \\ \sigma_e^{\text{eff}} \nabla \phi_e^2 + \sigma_m^{\text{eff}} \nabla \phi_m^2 - J_c \left(\phi_{\text{eq},c} - \frac{-\Delta H_{\text{orr}}}{2F} \right) - J_c \eta_c + S_T^{v-1} + S_T^{v-d} + J_{\text{cross}} \left(\phi_{\text{eq},a} - \frac{-\Delta H_{\text{hor}}}{2F} \right) + J_{\text{cross}} \eta_{\text{cross}} & \text{(cCL)} \end{cases} \quad (27)$$

where $\phi_{\text{eq},a}$ and $\phi_{\text{eq},c}$ represent the equilibrium potentials for HOR and ORR, respectively. ΔH_{hor} and ΔH_{orr} denote the enthalpy changes associated with HOR and ORR, respectively. Additionally, η_a , η_c and η_{cross} represent the overpotentials for HOR at the anode, ORR at the cathode, and HOR at the cathode, respectively. S_T^{v-1} and S_T^{v-d} represent the heat associated with phase changes of water, as detailed in our previous work. $J_{\text{cross}} \left(\phi_{\text{eq},a} - \frac{-\Delta H_{\text{hor}}}{2F} \right)$ is the reversible heat for HOR at the cathode, while $J_{\text{cross}} \eta_{\text{cross}}$ is the irreversible heat for HOR at the cathode.

This concludes the energy transport equation group. The resulting temperature field affects the hydrogen diffusion coefficient in Section 3.2.1 via the Arrhenius relation (Eq. (3)), closing the thermal-hydrogen coupling loop.

3.3 Reaction Kinetics

This section presents the electrochemical kinetics equation group (Eqs. (28)–(33)), defining reaction rates and overpotentials via Butler-Volmer equations for HOR and ORR.

As shown in Fig. 1c, the anode exclusively facilitates HOR, with reaction kinetics characterized by the Butler-Volmer equation [49]:

$$J_a = i_{0,a}^{\text{ref}} a_{\text{Pt}} \exp \left[-1400 \left(\frac{1}{T} - \frac{1}{353.15} \right) \right] \left(\frac{C_{H_2}}{C_{H_2}^{\text{ref}}} \right) \left[\exp \left(\frac{F}{RT} \eta_a \right) - \exp \left(-\frac{F}{RT} \eta_a \right) \right] \quad (28)$$

where i_0^{ref} is the reference exchange current density; a_{Pt} is the electrocatalytic surface area per unit volume; and C_i^{ref} denotes the reference molar concentration.

At the cathode, both ORR and HOR coexist due to H_2 permeation. The ORR rate is described by the Butler-Volmer equation, given as follows:

$$J_c = i_{0,c}^{\text{ref}} a_{\text{Pt}} \exp \left[-7900 \left(\frac{1}{T} - \frac{1}{353.15} \right) \right] \left(\frac{C_{O_2}}{C_{O_2}^{\text{ref}}} \right) \left[\exp \left(\frac{3F}{RT} \eta_c \right) - \exp \left(-\frac{F}{RT} \eta_c \right) \right] \quad (29)$$

The HOR rate at the cathode is sufficiently high due to the high electrode potential, and is solved using the following equation [23]:

$$J_{\text{cross}} = i_{0,a}^{\text{ref}} a_{\text{Pt}} \frac{C_{\text{H}_2,\text{diss}}}{C_{\text{H}_2}^{\text{ref}}} \exp \left[-1400 \left(\frac{1}{T} - \frac{1}{353.15} \right) \right] \left[\exp \left(\frac{F}{RT} \eta_{\text{cross}} \right) - \exp \left(-\frac{F}{RT} \eta_{\text{cross}} \right) \right] \quad (30)$$

where η_{cross} is the overpotential for the HOR inside the cCL.

The above-mentioned three overpotentials are solved using:

$$\eta_a = \phi_e - \phi_m \quad (31)$$

$$\eta_c = \phi_e - \phi_m - \phi_{\text{eq},c} \quad (32)$$

$$\eta_{\text{cross}} = \phi_e - \phi_m \quad (33)$$

where ϕ_e and ϕ_m are the electric potential and the ionic potential, respectively, ϕ_{eq} is the equilibrium potential.

This concludes the kinetics equation group. These current densities close the coupling loop by serving as source terms for the charge (Section 3.2.5), mass (Sections 3.2.1–3.2.3), and energy (Section 3.2.6) conservation equations.

3.4 Boundary Conditions

This section defines the boundary condition equation group (Eqs. (34)–(37)), specifying inlet stoichiometric flows, concentrations, and interfacial flux conditions.

The primary boundary conditions are as follows: the outlet pressures of both the anode and cathode are fixed to the respective backing pressures. The inlet flow rates of the gas mixtures on both the cathode and anode sides are determined by solving:

$$\mathbf{u}_a = \xi_a \frac{I(t)A_m}{2FA_{\text{ch}}C_{\text{H}_2}} \quad (34)$$

$$\mathbf{u}_c = \xi_c \frac{I(t)A_m}{4FA_{\text{ch}}C_{\text{O}_2}} \quad (35)$$

where ξ denotes the stoichiometric ratios of H_2 at the anode and O_2 at the cathode, $I(t)$ represents the time-varying current load, A_m denotes the area of the membrane-electrode assembly parallel to the flow channel direction, and A_{ch} represents the cross-sectional area of the flow channel. Additionally, C_{H_2} and C_{O_2} denote the concentrations of hydrogen and oxygen at the inlet, respectively, and are expressed as:

$$C_{\text{H}_2} = \frac{p_{g,a} - RH_a \times p_{\text{sat}}}{RT} \quad (36)$$

$$C_{\text{O}_2} = 0.21 \left(\frac{p_{g,c} - RH_c \times p_{\text{sat}}}{RT} \right) \quad (37)$$

where p_g , RH and p_{sat} are the gas pressures at the inlets, the relative humidities, and the saturation vapor pressure, respectively.

Furthermore, the pressure of the liquid mixture p_{ml} at the interface between the MPL and CL within the electrolyte is set to be in equilibrium with the liquid pressure. The H_2 crossover flux ($\mathbf{u}_{\text{H}_2,\text{diss}}C_{\text{H}_2,\text{diss}} - D_{\text{H}_2,\text{diss}}\nabla C_{\text{H}_2,\text{diss}}$) is maintained at zero at the MPL|CL interface. These boundary conditions drive the transient response of all preceding equation groups.

4 Results and Discussion

4.1 Model Implementation and Verification

The computational model was analyzed using COMSOL 6.1. Fig. 3a illustrates the outcomes of the independence test regarding the number of elements. This evaluation was performed under high current densities ranging from 1800 to 2000 mA cm⁻², reflecting the rigorous requirements for element quality at these elevated current levels. Notably, when the element count exceeds 14,000, the model demonstrates a nearly constant voltage prediction trend as the number of elements increases. To achieve an optimal balance between computational accuracy and time efficiency, the final selected number of elements is established at 14,000. Fig. 3b,c presents a comparison of model-predicted data with experimentally measured data. The experimental conditions were meticulously aligned with the simulation parameters, which included a temperature of 90°C, backing pressures of 130 kPa at the anode and 120 kPa at the cathode, stoichiometric ratios of 1.7 at the anode and 3.0 at the cathode, and a relative humidity of 0.5 at both the anode and cathode. Additionally, the membrane thickness was set at 8 μm. Fig. 3b illustrates the correlation between the steady-state polarization curves obtained from the 3D model and the experimental data, demonstrating a maximum error of 2%. In Fig. 3c, a significant correlation is observed in the transient cell voltage response under a step load condition, with a current step from 600 to 1600 mA cm⁻² applied between 2 and 3 s, with an error margin also below 2%. Both the simulations and experiments effectively capture the undershoot voltage of approximately 0.01 V, which can be attributed to instantaneous local oxygen starvation and low levels of membrane hydration. This validation relies on global performance metrics (steady-state polarization and transient voltage), which is widely adopted in the literature to establish model credibility before extending predictions to local phenomena [27,50].

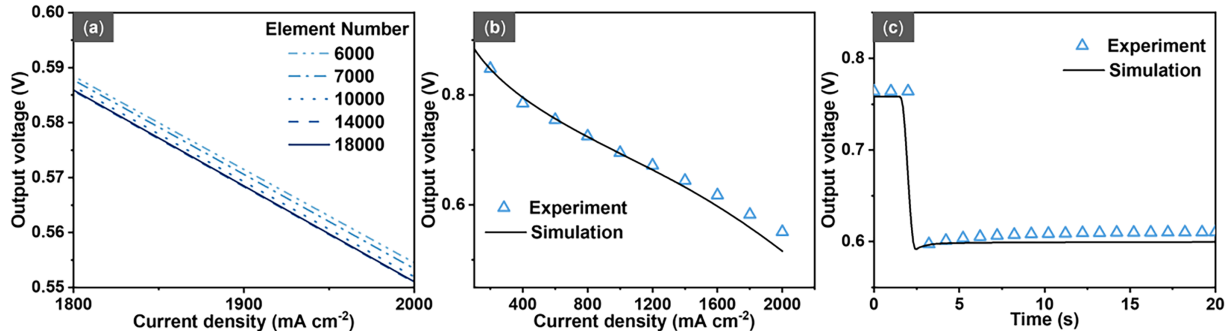


Figure 3: Model validation under 90°C, 130/120 kPa (anode/cathode), 50% RH, stoichiometric ratios of 1.7/3.0 with 8 μm membrane: (a) Mesh independence test at 1800–2000 mA cm⁻² showing convergence at 14,000 elements; (b) Steady-state polarization curve comparison between simulation and experiment (maximum error of 2%); (c) Transient voltage response under step load from 600 to 1600 mA cm⁻² from 2 to 3 s over 1 s, showing ~0.01 V undershoot (error below 2%).

Fig. 4a presents a comparative analysis of the model results with and without considering the impact of H₂ crossover on fuel cell performance, mass transfer, and heat transfer. The findings indicate that neglecting H₂ crossover leads to a substantial overestimation of the output voltage, particularly the open circuit voltage (OCV). Specifically, the OCV is 0.972 V when H₂ crossover effects are accounted for, while it rises to 1.182 V when these effects are ignored, resulting in a difference of 210 mV. This discrepancy is attributed to the formation of localized H₂-O₂ fuel cells, aligning with the simulation results reported by Vilekar and Datta [23]. However, the influence of these localized fuel cells diminishes notably with increasing load, as their effect on electrode potential decreases [22]. Additionally, when H₂ crossover effects are not considered,

the cell average temperature is underestimated by approximately 0.2°C across the entire current density range. This increase in temperature observed with H_2 crossover is primarily due to the additional heat generated during the HOR and ORR at the cathode [50]. Fig. 4b further presents a comparison between the H_2 crossover current density predicted by the model and the experimental data under identical operating conditions: a Nafion 112 membrane with a thickness of $50\ \mu\text{m}$, at 80°C , with saturated gases, and at OCV. The model predictions demonstrate excellent agreement with the experimental data, illustrating a consistent trend where the H_2 crossover current increases with the rise in H_2 partial pressure on the anode side. This correlation emphasizes the accuracy of the model in capturing the effects of H_2 crossover under specified conditions.

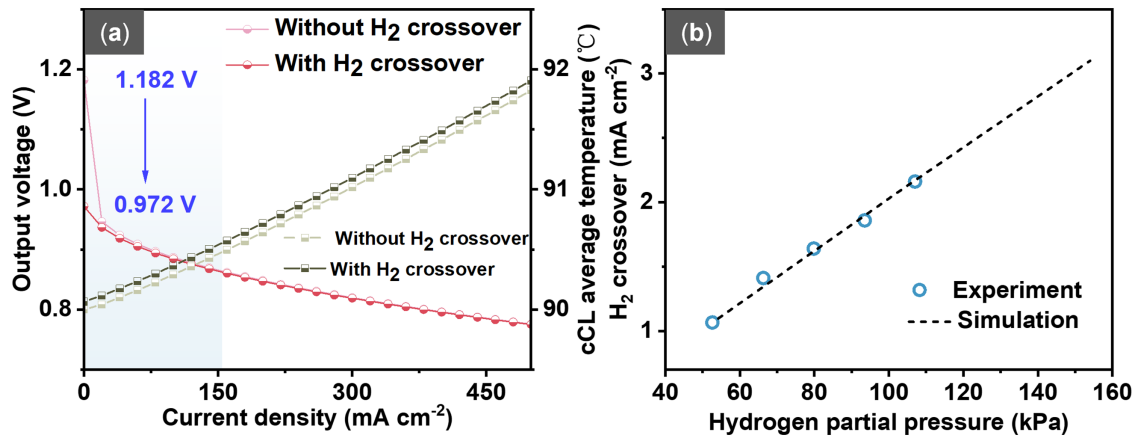


Figure 4: Impact of H_2 crossover on fuel cell performance and model validation: (a) Comparison of output voltage and cCL average temperature with and without H_2 crossover (OCV reduction of 210 mV from 1.182 to 0.972 V, and $\sim 0.2^{\circ}\text{C}$ temperature increase when crossover is considered); (b) Validation of H_2 crossover current density vs. anode H_2 partial pressure under Nafion 112, $50\ \mu\text{m}$ thickness, 80°C , saturated gases, OCV conditions (simulation vs. experimental data [15]).

4.2 Effect of Membrane Thickness

The influence of membrane thickness on fuel cell performance and the concentration of dissolved H_2 under dynamic load conditions is illustrated in Fig. 5. Specifically, Fig. 5a demonstrates a significant reduction in the output voltage of the fuel cell as the current load increases rapidly from 10 to $1800\ \text{mA cm}^{-2}$ between 9.9 and 10.1 s (within 0.2 s), which is accompanied by a notable voltage undershoot. Fig. 5b shows that after 30 s, the voltage decreases from 0.59 to 0.55, 0.52, and 0.49 V as the thickness of the PEM is increased from 5 to 10, 15, and $20\ \mu\text{m}$, respectively. Correspondingly, the voltage undershoots, defined as the difference between the voltage at 30 s and the minimum value observed at 10 s, increases to 0.013, 0.024, and 0.036 V from a baseline of 0.004 V, respectively.

These results suggest a heightened power loss during both steady-state and dynamic operations, which can be attributed to the increased proton-conducting resistance associated with thicker membranes. Furthermore, it is critical to note that the thickness of the PEM not only impacts the performance of the fuel cell but also induces variations in internal H_2 crossover.

Fig. 5c illustrates the dynamic responses of dissolved H_2 concentration within the PEM, and shows that the steady-state dissolved H_2 concentration exhibits a positive correlation with the thickness of the PEM, with a thicker membrane facilitating higher concentrations. Additionally, an undershoot phenomenon is observed for dissolved H_2 concentration when the load undergoes rapid increases, with slight variations depending on the PEM thickness, as depicted in Fig. 5d.

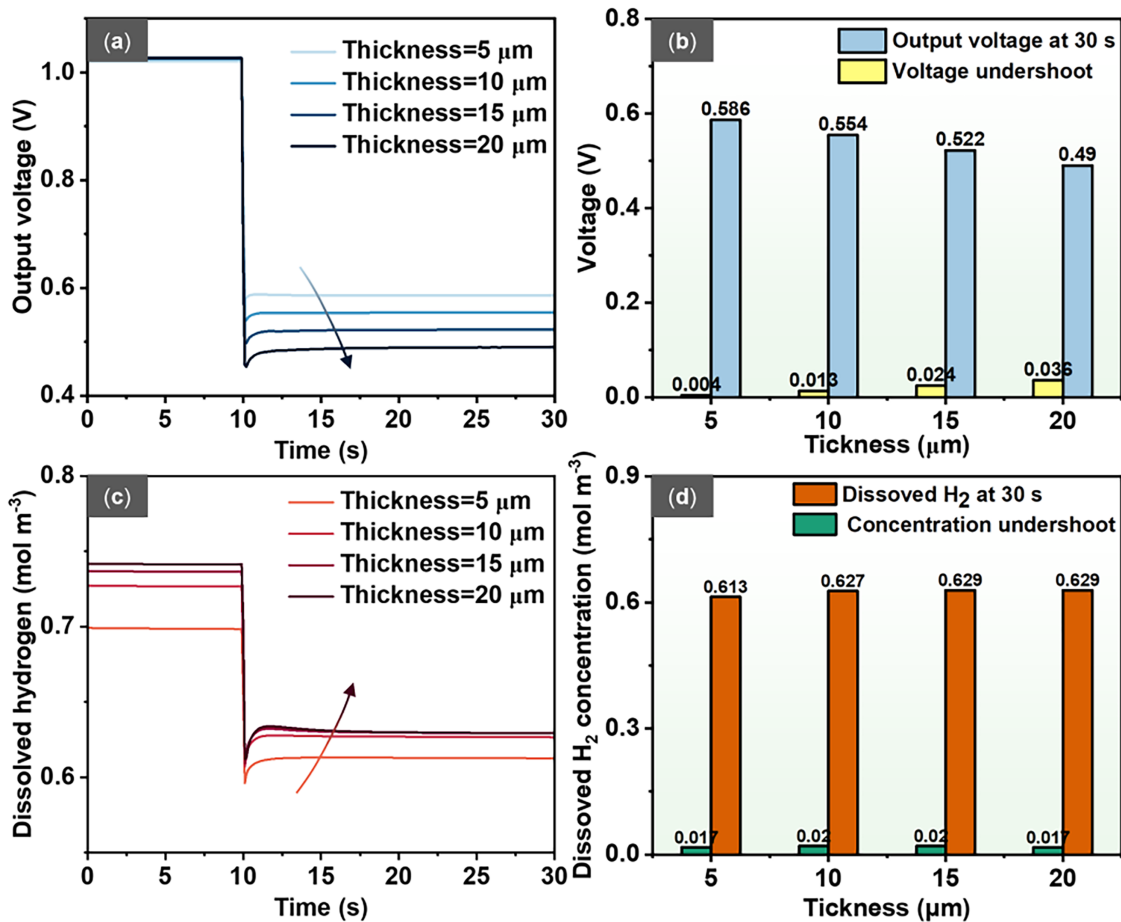


Figure 5: Effect of membrane thickness on fuel cell dynamic performance under a step load change from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 s to 10.1 s) with total simulation time of 30 s, operating conditions: 90°C, 130/120 kPa (anode/cathode), 50% RH: (a) Dynamic response of output voltage from 9.9 s to 10.1 s for PEM thicknesses of 5, 10, 15, and 20 μm; (b) Variations of voltage undershoot and steady-state voltage at 30 s with PEM thickness; (c) Dynamic responses of dissolved H₂ concentration at the anode-PEM interface for PEM thicknesses of 5, 10, 15, and 20 μm; (d) Variations of dissolved H₂ concentration undershoot and steady-state concentration at 30 s with PEM thickness.

The undershoot in dissolved H₂ concentration can be attributed to the dynamic interplay of local H₂ gas concentration, temperature, and membrane water content. Fig. 6a clearly demonstrates that as the load increases, there is a rapid decrease in H₂ gas concentration within the aCL, leading to a significant decline in the dissolved H₂ concentration, as observed in Fig. 5c. This decline is governed by the proportional relationship between H₂ gas partial pressure and the saturated dissolved H₂ concentration, as described by Eq. (7). However, over time, the dissolved H₂ concentration gradually increases, eventually reaching a new steady state, which gives rise to the undershoot phenomenon. Notably, the H₂ gas concentration within the aCL remains relatively constant during this process. The gradual increase in dissolved H₂ can be attributed to enhanced H₂ crossover resulting from increased local membrane water content and temperature, which facilitate H₂ dissolution and permeation through the membrane.

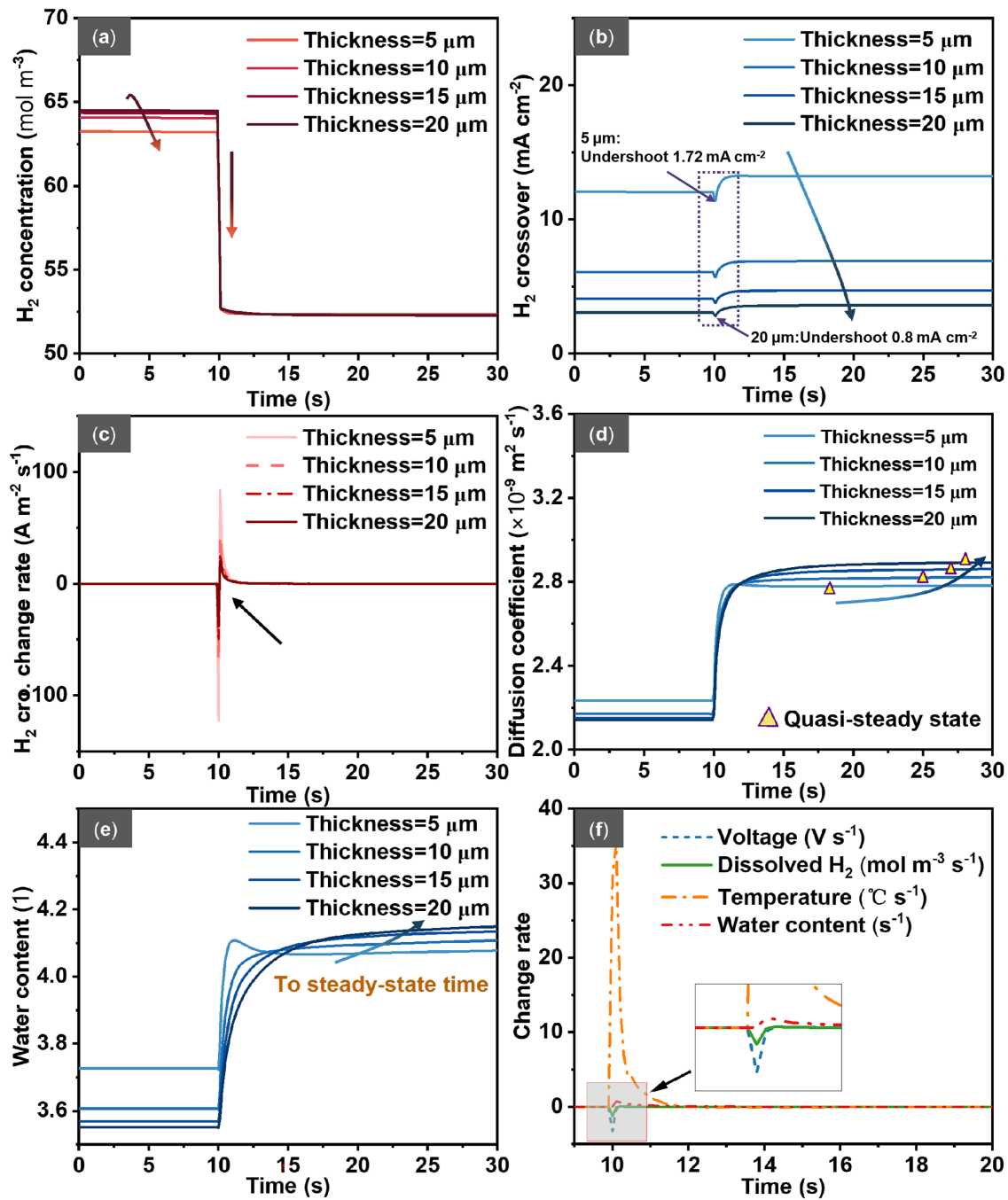


Figure 6: Dynamic response mechanisms of hydrogen transport and membrane properties under a step load from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 s to 10.1 s) at 90°C, 130/120 kPa (anode/cathode), 50% RH for various PEM thicknesses (5–20 μm): (a) Average H₂ gas concentration in the anode catalyst layer; (b) H₂ crossover current density through the PEM; transient minima and post-step recovered values are approximately 2.8 and 3.6 mA cm⁻² for the 20 μm membrane, and 11.4 and 13.2 mA cm⁻² for the 5 μm membrane, respectively; the undershoot, defined as the difference between the post-step recovered value and the transient minimum, ranges from approximately 0.8 mA cm⁻² for the 20 μm membrane to 1.72 mA cm⁻² for the 5 μm membrane; (c) Change rate of H₂ crossover current density; (d) Diffusion coefficient of dissolved H₂; (e) Water content in the PEM; (f) Comparative change rates of fuel cell voltage, dissolved H₂ inside PEM, PEM temperature, and PEM water content illustrating the temporal sequence of transient responses.

As shown in Fig. 6b, the H₂ crossover current density through the PEM demonstrates an almost fourfold reduction as the PEM thickness increases from 5 to 20 μm. For the 5 μm membrane, the H₂ crossover current density decreases to a transient minimum of approximately 11.4 mA cm⁻² and subsequently recovers to approximately 13.2 mA cm⁻². For the 20 μm membrane, the corresponding transient minimum and post-step recovered value are approximately 2.8 and 3.6 mA cm⁻², respectively. This substantial decrease can be ascribed to the longer transport pathways and the greater transport resistance inherent to thicker membranes, which hinder the transfer of dissolved H₂ from the aCL to the cathode. Moreover, this phenomenon also results in a decrease in the rate of change of H₂ crossover with increasing thickness of the PEM. From the perspective of diffusion mechanisms regulating H₂ crossover, a thicker membrane creates a longer diffusion path, which consequently delays the changes in H₂ crossover [28,33,51]. Additionally, when examining the rate of change of H₂ crossover in relation to varying loads, Fig. 6c illustrates that the change rate initially follows a negative trajectory and then shifts to a positive trend as the load increases. This pattern is consistent with the H₂ crossover current density presented in Fig. 6b, where an initial decline is observed, followed by a notable increase around the 10-s mark. The resulting undershoot, defined as the difference between the post-step recovered value and the transient minimum, ranges from approximately 0.8 mA cm⁻² for the 20 μm membrane to 1.72 mA cm⁻² for the 5 μm membrane.

The initial decrease can be attributed to a diminishing gradient of dissolved H₂ from the anode to the cathode. Conversely, the subsequent increase is linked to an enhanced diffusion coefficient, which arises from the hydration of the membrane, as illustrated in Fig. 6d. The increase in PEM thickness leads to a longer time for the diffusion coefficient in Fig. 6d to reach a new steady state. This delay can be attributed to the slower evolution of membrane water, as observed in Fig. 6e. The rates of change of voltage, dissolved H₂, local temperature, and water content are summarized in Fig. 6f. The results indicate that the output voltage exhibits the fastest response, followed by dissolved H₂, membrane water, and temperature. This suggests that the dynamic response of dissolved H₂ is faster compared to heat and water transfer.

Fig. 7 presents the 3D distributions of dissolved H₂, H₂ crossover, and water content within a 15 μm PEM during load steps. As the load increases from 9.9 s to 10 s, the dissolved H₂ near the aCL decreases noticeably, which is attributed to the rapid decline in H₂ gas concentration inside the aCL. This decrease in dissolved H₂ concentration gradient from the anode to the cathode corresponds to a reduction in H₂ crossover flux from 4.0–4.1 to 3.5–4.0 mA cm⁻², as depicted in Fig. 7e,f. These findings are consistent with the average results shown in Figs. 5c and 6b. Between 10.1 s and 11 s, an increase in dissolved H₂ is observed in the inlet gas area adjacent to the anode side. This increase can be attributed to the elevated local temperature and the enhanced H₂ crossover flux, both of which facilitate a greater influx of H₂ gas from the aCL into the PEM. Over the same period, a significant increase in membrane water content is observed beneath the rib region, which primarily explains the rise in H₂ crossover from approximately 3.5–4.0 to 4.1–4.9 mA cm⁻². It is important to note that the expansion of Nafion clusters induced by water infiltration has previously been reported to enhance gas permeation [10,52,53].

At t = 10 s and 10.1 s, an unexpected increase in H₂ crossover is observed under the channel region, with values 0.1–0.7 mA cm⁻² higher than those in the rib region, which contrasts with the trend in membrane water distribution. This variation can be attributed to the impact of the concentration gradient on H₂ crossover. During these time intervals, both the channel and rib regions exhibit low H₂ diffusion coefficient due to membrane dehydration. However, the intensified H₂ concentration gradient leads to an increased occurrence of H₂ crossover specifically in the channel region. As depicted in Fig. 7b,c, the dissolved H₂ concentration is notably higher under the channel region on the anode side, resulting in a localized concentration disparity that promotes a higher incidence of H₂ crossover in the channel region.

In contrast, at 11 s, the substantial increase in membrane water content beneath the rib region (Fig. 7l) correlates with a 0.1–0.4 mA cm⁻² higher H₂ crossover relative to the channel region. This association is captured by the water-dependent diffusion coefficient (Eq. (2)) in the present continuum-scale formulation. While enhanced permeation in these hydrated regions is speculated based on the model results, it should be noted that the present macroscopic model does not explicitly resolve nanoscale structural features such as distinct water channel networks or ionic cluster morphology, consistent with the crossover distribution shown in Fig. 7h.

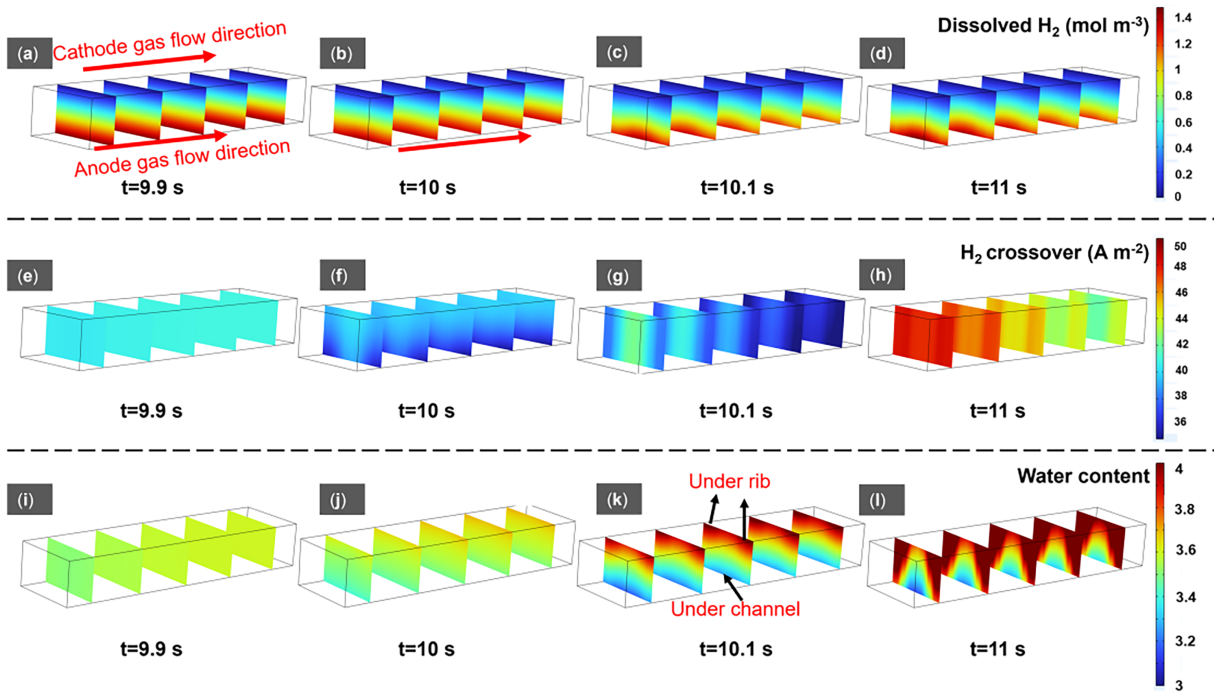


Figure 7: Transient spatial distributions within a 15 μm PEM during a step load change from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 s to 10.1 s) at 90°C, 130/120 kPa (anode/cathode), and 50% RH, shown at critical moments ($t = 9.9, 10, 10.1,$ and 11 s): (a–d) Dissolved H₂ concentration; (e–h) H₂ crossover flux; (i–l) Water content, illustrating the contrasting evolution between channel and rib regions during the load transient.

4.3 Effect of Stoichiometric Ratio

In this study, we investigate the impact of the anode stoichiometric ratio on H₂ crossover in a PEM fuel cell. The results presented in Fig. 8a demonstrate a significant increase in dissolved H₂ concentration within the PEM at higher anode stoichiometric ratios, particularly under high current loads. For instance, at a current density of 1800 mA cm⁻² and 30 s, the dissolved H₂ concentrations are measured at 0.56, 0.60, 0.63, and 0.65 mol m⁻³ for H₂ stoichiometric ratios of 0.9, 1.3, 1.7, and 2, respectively (Fig. 8b).

This increase can be attributed to the elevated H₂ gas concentration within the aCL, as further supported by the data shown in Fig. S1. The amplified dissolved H₂ concentration further increases H₂ crossover flux and its dynamic undershoot. As shown in Fig. 8c, the H₂ crossover flux values at 10.1 s are 3.4, 3.63, 3.78, and 3.85 mA cm⁻² for H₂ stoichiometric ratios of 0.9, 1.3, 1.7, and 2, respectively. Following the load change at 30 s, the increased dissolved H₂ at anode with stoichiometric ratio increases H₂ crossover with values of 4.27, 4.57, 4.73, and 4.81 mA cm⁻², respectively. This results in a more pronounced undershoot of H₂ crossover flux, with values of 0.87, 0.94, 0.95, and 0.96 mA cm⁻², respectively.

Fig. 8d further compares the change rates of H₂ crossover. As the anode stoichiometric ratio increases, the initial rate of decrease in H₂ crossover becomes less pronounced, while the subsequent rate of increase is markedly amplified. These trends can be attributed to the enhanced dissolved H₂ gradient between the anode and cathode. This increased gradient mitigates the transient decline in H₂ crossover resulting from the decrease in H₂ concentration, while simultaneously promoting the subsequent increase due to improved membrane hydration.

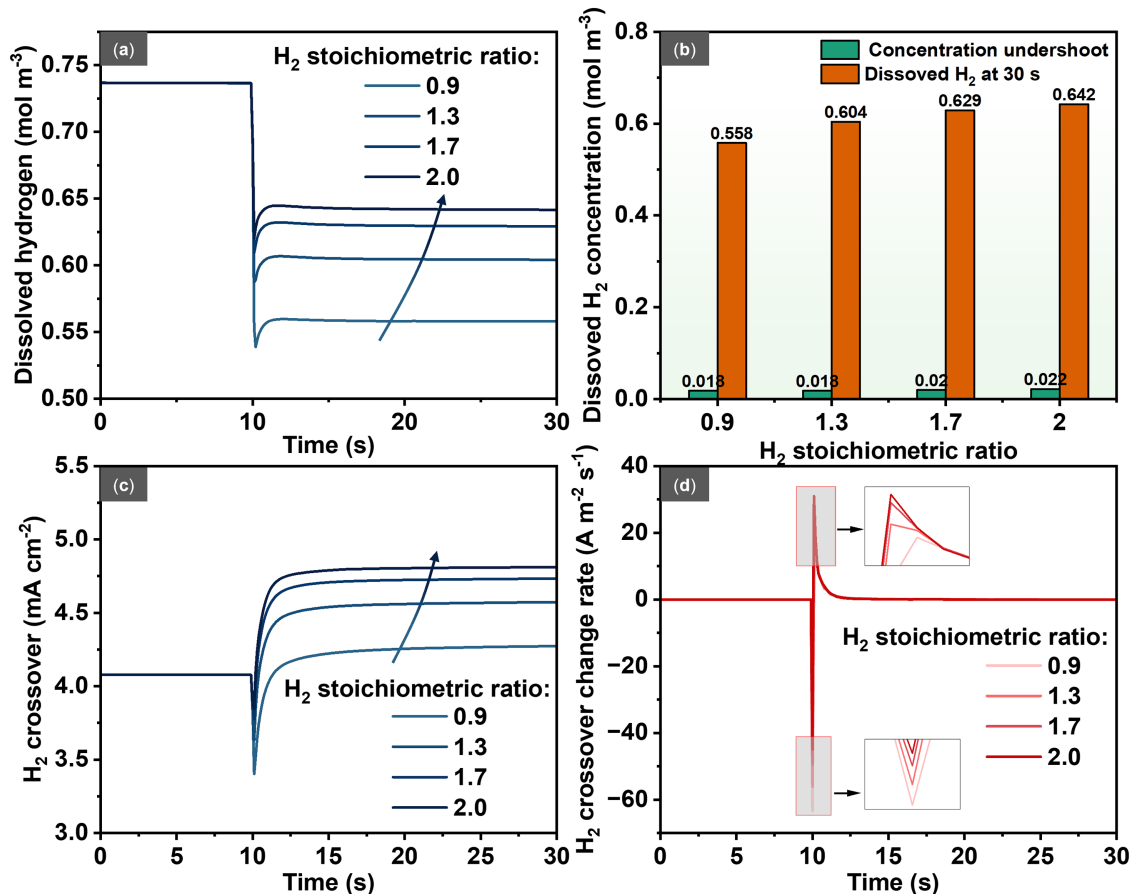


Figure 8: Effect of anode stoichiometric ratio ($\xi_a = 0.9, 1.3, 1.7, 2$) on hydrogen transport dynamics under a step load from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 s to 10.1 s) at 90°C, 130/120 kPa (anode/cathode), and 50% RH: (a) Dynamic responses of dissolved H₂ concentration; (b) Dissolved H₂ concentration undershoot and steady-state values at 30 s; (c) Dynamic responses of H₂ crossover flux (values of 3.4–3.85 mA cm⁻² at 10.1 s and 4.27–4.81 mA cm⁻² at 30 s); (d) Change rates of H₂ crossover illustrating the transient undershoot phenomenon (0.87–0.96 mA cm⁻²).

Fig. 9a–d illustrates the distribution of dissolved H₂ within the PEM at 10.1 s. The concentration of dissolved H₂ increases with the anode stoichiometric ratio, peaking in the region near the inlet on the anode side. This elevated concentration contributes to a larger area of high H₂ crossover within the PEM, as shown in Fig. 9e–h. The H₂ crossover flux varies across the anode stoichiometric ratios, ranging from 2.7–4.1 mA cm⁻² for 0.9, 3.2–4.2 mA cm⁻² for 1.3, 3.4–4.2 mA cm⁻² for 1.7, and 3.5–4.3 mA cm⁻² for 2. Notably, the channel region exhibits a more pronounced H₂ crossover due to a steeper H₂ concentration gradient. Additionally, H₂ crossover diminishes along with the gas flow direction, which can be attributed to the decrease in anode H₂ gas resulting from consumption during the electrochemical reaction.

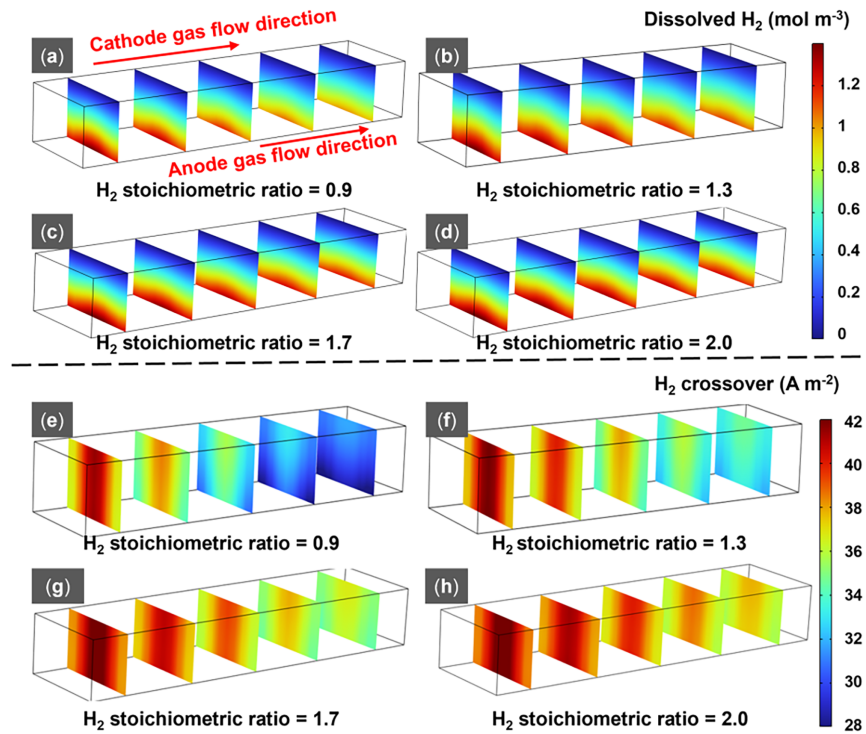


Figure 9: Spatial distributions of dissolved H_2 (a–d) and H_2 crossover flux (e–h) within the $15\ \mu\text{m}$ PEM at $t = 10.1\ \text{s}$ as the current load increases from 10 to $1800\ \text{mA cm}^{-2}$ at 90°C , 130/120 kPa, and 50% RH for anode stoichiometric ratios of $\xi_a = 0.9, 1.3, 1.7,$ and 2 , illustrating the higher crossover in the channel region and decreasing trend along the flow direction.

4.4 Effect of Inlet Gas Humidity

In Fig. 10a, higher anode humidity levels result in a significant increase in output voltage under high loads. At 30 s, the voltage values are 0.448, 0.488, 0.522, and 0.563 V at humidity levels of 20%, 30%, 50%, and 80%, respectively (Fig. 10b). The magnitude of voltage undershoot decreases accordingly, with values of 0.142, 0.073, 0.023, and $-0.02\ \text{V}$, respectively.

These changes can be attributed to a reduction in ohmic polarization loss due to membrane hydration.

Fig. 10c summarizes the dynamic response of dissolved H_2 within the PEM, showing a significant decrease with increasing anode humidity, regardless of the load condition. At 30 s, the dissolved H_2 concentration is 0.703, 0.678, 0.629, and $0.556\ \text{mol m}^{-3}$ for anode humidities of 20%, 30%, 50%, and 80%, respectively (Fig. 10d).

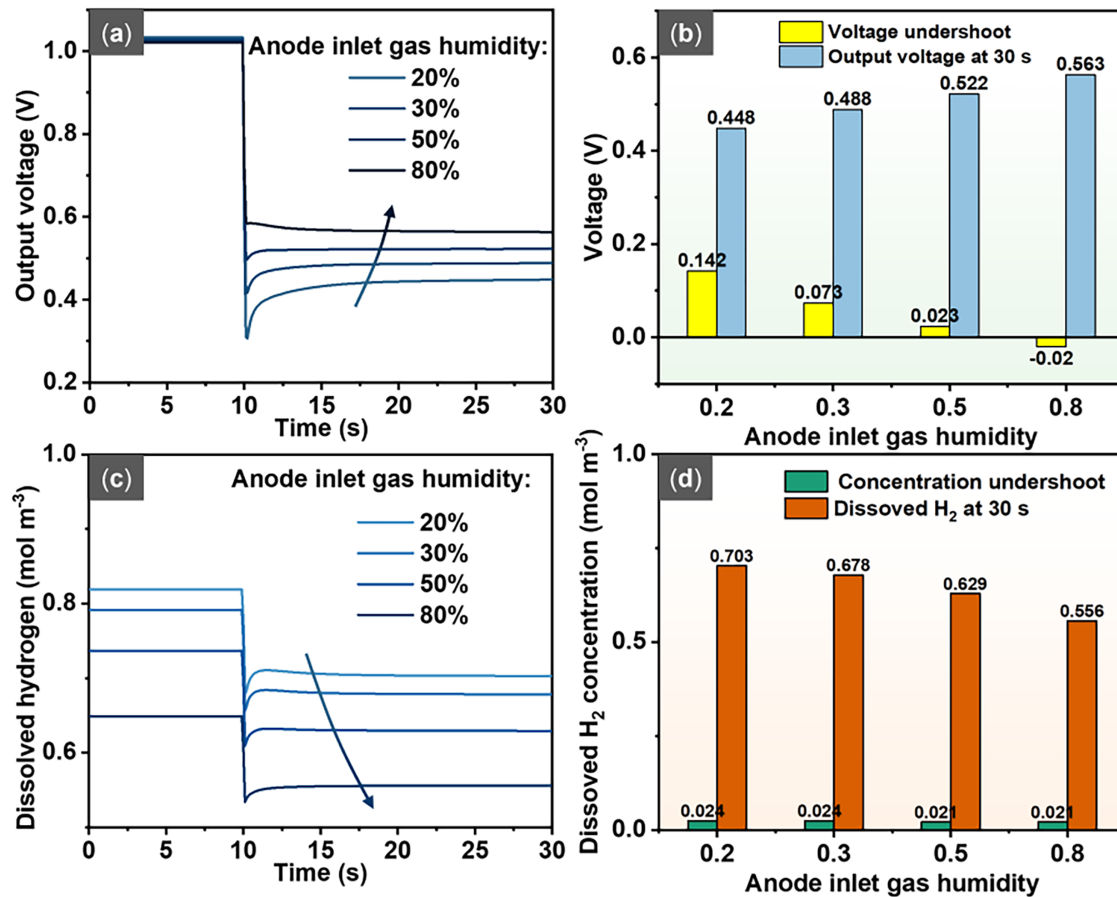


Figure 10: Effect of anode inlet humidity (20%, 30%, 50%, 80% RH) on fuel cell performance under a step load from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 s to 10.1 s) at 90°C, 130/120 kPa for the 15 μm PEM: (a) Dynamic response of output voltage (increasing from 0.448 to 0.563 V at 30 s with rising anode humidity); (b) Voltage undershoot and steady-state voltage at 30 s (undershoot decreasing from 0.142 to -0.02 V with rising anode humidity); (c) Dynamic response of dissolved H₂ concentration (decreasing from 0.703 to 0.556 mol m⁻³ at 30 s with rising anode humidity); (d) Dissolved H₂ undershoot and steady-state concentration at 30 s.

This decrease is a result of increased water vapor diluting the H₂ gas concentration within the aCL (refer to Fig. 11a), leading to a subsequent decrease in dissolved H₂ concentration.

Fig. 11b summarizes the dynamic responses of H₂ crossover flux, revealing a significant increase from 3.5–4.6 to 5.1–5.6 mA cm⁻² as anode humidity rises from 20% to 80%. This increase is attributed to the enhanced gas diffusion coefficient within the PEM (Fig. 11c) due to the increased water content in the membrane. Fig. 11d compares the change rates of H₂ crossover under different anode gas humidities. The rate of decrease in H₂ crossover demonstrates a significant increase with rising humidity, while the subsequent rate of increase exhibits a slight decrease. These variations are closely associated with the dynamics of local H₂ gas concentration and membrane water content. When the load abruptly increases, the concentration of anode H₂ gas decreases significantly with rising humidity, resulting in a more rapid decline in H₂ crossover. However, the subsequent increase in membrane water content diminishes with higher humidity (Fig. S2), leading to a slower rebound in H₂ crossover at elevated humidity levels.

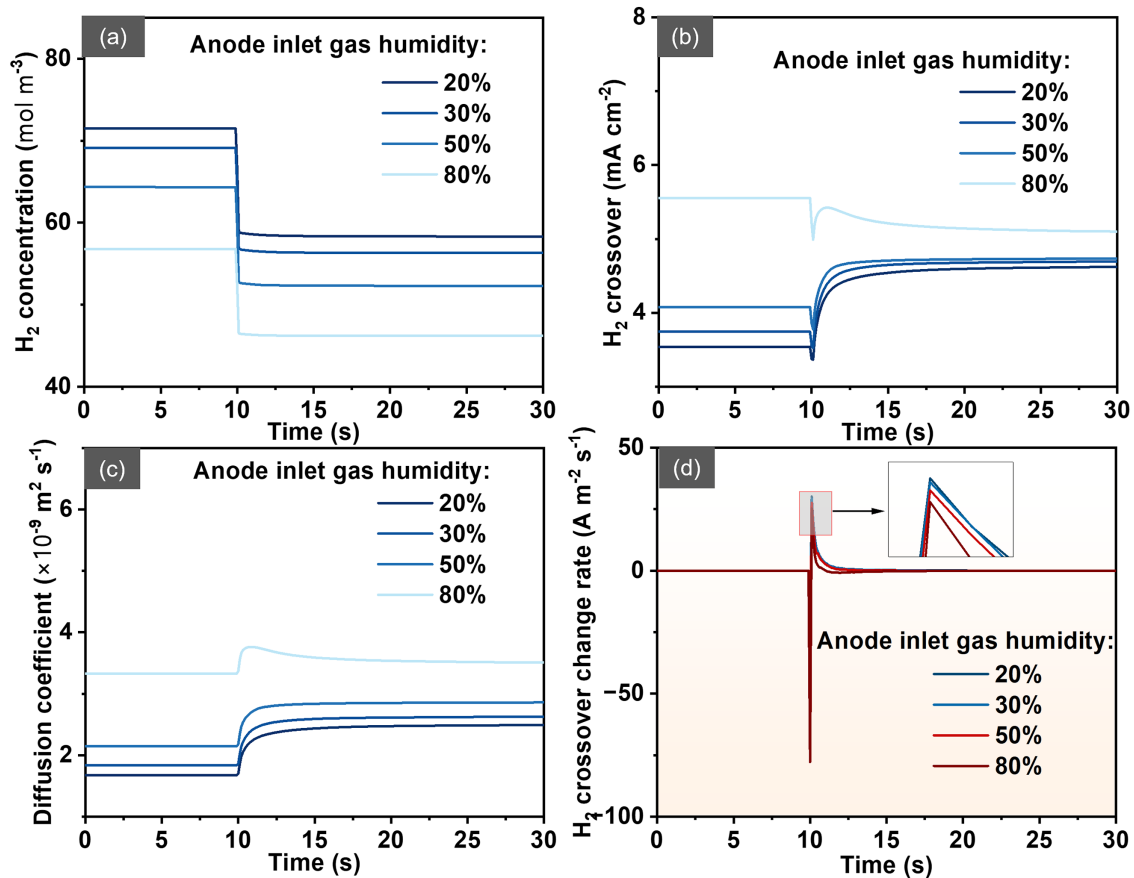


Figure 11: Mechanistic analysis of hydrogen transport under a step load from 10 to 1800 mA cm⁻² within 0.2 s (from 9.9 to 10.1 s) at 90°C, 130/120 kPa for the 15 μm PEM with anode humidity levels (20%–80% RH), illustrating the competition between dilution and diffusion enhancement: (a) H₂ gas concentration in the aCL (decreasing with humidity due to dilution); (b) H₂ crossover flux (increasing from 3.5–4.6 to 5.1–5.6 mA cm⁻² despite dilution); (c) Diffusion coefficient of dissolved H₂ (increasing with membrane hydration); (d) Change rates of crossover flux.

Fig. 12a–d summarizes the distributions of H₂ crossover within the PEM at 11 s. Consistent with the average dynamic response, H₂ crossover increases by approximately 0.8 mA cm⁻² throughout the PEM as humidity increases, particularly in the region near the gas inlet under the rib. This increase in H₂ crossover with humidity is primarily attributed to the rise in membrane water content. Fig. 12e–h demonstrates a significant increase in membrane water content as anode humidity increases, especially in the area under the rib. However, in contrast to the H₂ crossover distribution, the water content near the outlet is higher due to product accumulation. This discrepancy can be explained by the more pronounced H₂ concentration gradient between the anode and cathode, which leads to more severe H₂ crossover in the inlet area.

4.5 Design Implications and Diagnostic Insights

The parametric study yields several quantitative guidelines for automotive PEMFC operation. First, membrane thickness should be optimized between 10 and 15 μm: a 5 μm membrane increases H₂ crossover fourfold (11.4–13.2 mA cm⁻²) posing safety concerns, whereas a 20 μm membrane causes higher ohmic loss (voltage drop ~0.1 V at 1800 mA cm⁻²). Second, under dynamic loading, an anode stoichiometric ratio of 1.3–1.7 is recommended: $\xi_a = 0.9$ yields the smallest crossover undershoot (0.87 mA cm⁻²) but may compromise

reactant distribution along the channel, while $\xi_a = 2.0$ gives a larger undershoot (0.96 mA cm^{-2}) and higher steady-state crossover (4.81 mA cm^{-2}); the intermediate range balances dynamic response and fuel efficiency. Inlet humidity of 50%–80% is preferred: 80% raises steady-state crossover to $5.1\text{--}5.6 \text{ mA cm}^{-2}$ but improves voltage by 115 mV at 1800 mA cm^{-2} . For diagnostics, the predicted open-circuit voltage drop of 210 mV and local cathode temperature rise of 0.2°C can serve as non-invasive indicators of membrane integrity. Moreover, the sign change of dJ_{cross}/dt (from negative to positive) during load steps, quantified in Figs. 6c, 8d and 11d, could be used in real-time control algorithms to estimate membrane hydration state.

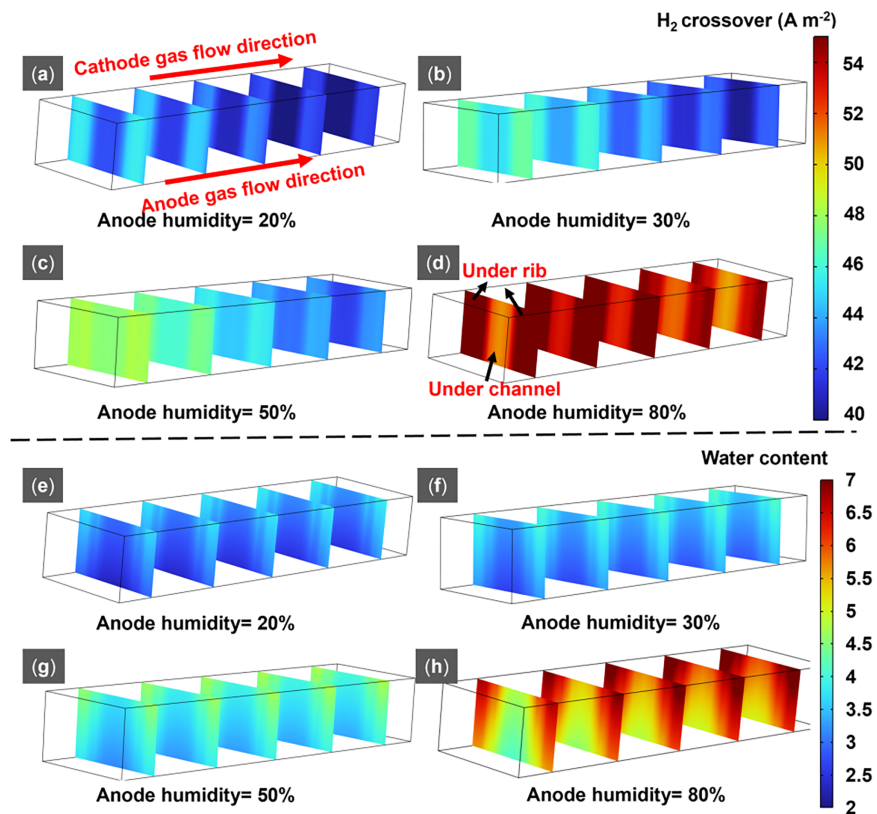


Figure 12: Spatial distributions within the $15 \mu\text{m}$ PEM at $t = 11 \text{ s}$ under step load conditions ($10 \rightarrow 1800 \text{ mA cm}^{-2}$) at 90°C , 130/120 kPa for anode humidity levels (20%, 30%, 50%, 80% RH): (a–d) H_2 crossover flux (increasing by $\sim 0.8 \text{ mA cm}^{-2}$, peak values near inlet under rib); (e–h) Water content (higher near outlet due to product accumulation).

5 Conclusions

This study presents a novel three-dimensional model designed to investigate H_2 crossover phenomena in PEMFCs under dynamic loading conditions. The model incorporates essential processes including hydrogen dissolution, diffusion, convection, release, and electrochemical reactions. A particular emphasis is placed on the significant impact of H_2 crossover on both the open circuit voltage and the local thermal dynamics within the fuel cell. Using the developed model, this research systematically explores the effects of dynamic loading profiles, membrane thickness, stoichiometric ratios, and inlet humidity on H_2 crossover evolution along a single channel. The numerical simulations reveal complex interactions among these parameters and their influence on H_2 crossover dynamics. The mechanisms underlying these observations are rigorously examined. The principal findings are summarized as follows:

- (1) **Spatiotemporal Variations of H₂ Crossover:** H₂ crossover within the PEM exhibits significant spatiotemporal variation influenced by the current load, membrane hydration level, and anode H₂ gas concentration. Under low current loads, the channel region experiences higher H₂ crossover levels, exceeding those in the rib region by 0.1–0.7 mA cm⁻², primarily due to a steeper H₂ concentration gradient. Conversely, at high current loads, H₂ crossover in the rib region increases markedly, surpassing that in the channel region by 0.1–0.4 mA cm⁻²; this phenomenon is linked to membrane swelling resulting from enhanced hydration.
- (2) **Impact of Membrane Thickness:** Reducing the PEM thickness from 20 to 5 μm results in a decrease in the overall power loss of the fuel cell but significantly increases H₂ crossover, increasing approximately fourfold, from 2.8–3.6 to 11.4–13.2 mA cm⁻². This heightened crossover is primarily attributed to the shorter transport path, which leads to enhanced levels of dissolved H₂ and dynamic undershoot effects. Notably, the rate of change in H₂ crossover during load variations increases as membrane thickness decreases, demonstrating an inverse relationship between diffusion and membrane thickness.
- (3) **Effect of Stoichiometric Ratio:** Increasing the anode stoichiometric ratio from 0.9 to 2.0 markedly enhances the dissolved H₂ gradient at elevated current loads, resulting in an increase in H₂ crossover from approximately 2.7–4.1 to 3.5–4.3 mA cm⁻², particularly in the channel region. As the load rapidly increases, the H₂ crossover flux initially declines sharply, followed by a gradual rise and eventual stabilization, creating an undershoot. The initial decline is due to the reduction in anode H₂ concentration, while the subsequent rise is driven by increased membrane water content. An increased stoichiometric ratio raises the anode H₂ concentration, which slows the initial decline rate and accelerates the rate of the subsequent increase.
- (4) **Influence of Anode Humidity:** Enhancing anode humidity from 20% to 80% improves overall fuel cell performance by 115 mV at 1800 mA cm⁻² and decreases voltage undershoot by 142 mV due to membrane hydration. It also concurrently leads to a decrease in dissolved H₂ due to the dilution effect of water vapor. However, H₂ crossover flux increases from 3.5–4.6 to 5.1–5.6 mA cm⁻², particularly near the gas inlet under the rib, which is attributed to an improved gas permeation coefficient resulting from membrane hydration. Furthermore, the interplay between local H₂ concentration and membrane water content significantly amplifies the rate of decrease in H₂ crossover at the moment of load change.
- (5) **Model Limitations:** The present model has several simplifications that should be acknowledged:
 - (i) Single-channel geometry, which does not capture multi-channel or stack-level interactions;
 - (ii) Validation based on global metrics (e.g., polarization curves, total crossover current) rather than spatially resolved local measurements (e.g., PCB segmented cell techniques)—direct local validation would be ideal but is beyond the current scope, a limitation not uncommon in numerical modeling studies of this nature;
 - (iii) Long-term degradation evolution (e.g., chemical aging, catalyst dissolution) is not resolved in the current transient framework;
 - (iv) The membrane is assumed defect-free, neglecting potential pinholes or inhomogeneities that may exist in real membranes;
 - (v) Although the model is non-isothermal, the detailed effects of local heating (e.g., hot spots and their impact on mechanical stress) are not quantitatively analyzed;
 - (vi) Mechanical degradation phenomena (e.g., membrane swelling, creep, compressive stress relaxation) relevant to safety and lifetime are not considered;
 - (vii) Parametric sensitivity of key transport parameters (H₂ diffusion coefficient, dissolution rate, and hydraulic permeability) within ±20%–50% uncertainty bounds has not been quantified;

- (viii) Characteristic time scales (τ) for gas-phase H₂ transport, membrane hydration, and thermal diffusion, and their explicit correlation with the sign change of H₂ crossover flux, have not been quantitatively decomposed.

These limitations do not undermine the main mechanistic conclusions but highlight priorities for future model extensions.

Finally, this study provides an advanced theoretical model for investigating H₂ crossover dynamics. The simulation results significantly enhance our understanding of transient hydrogen transport in PEMFCs. Future work will focus on: (i) experimental validation of the predicted local distributions and extension to multi-cell stack configurations; (ii) comprehensive parametric sensitivity analysis of transport parameters within $\pm 20\%$ – 50% bounds; (iii) quantitative decomposition of characteristic time scales via linearized perturbation analysis; and (iv) experimental validation of H₂ crossover under non-OCV (load) conditions, thereby bridging the gap between the present modeling framework and practical fuel cell operation.

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